

The Imprint of The Extragalactic Background Light in the Gamma-Ray Spectra of Blazars

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Abstract: The light emitted by stars and accreting compact objects through the history of the universe is encoded in the intensity of the Extragalactic Background Light (EBL). Knowledge of the EBL is important to understand the nature of star formation and galaxy evolution. Direct measurements of the EBL intensity are difficult due to strong foreground emissions within our Galaxy and Solar system. However, the EBL density can be measured indirectly due to the absorption of gamma rays on this photon field. The absence of such absorption features in the spectra of individual gamma-ray sources had been used in the past to set limits on the level of the EBL. Although this feature is indeed difficult to constrain for a single source, it could be detected collectively in the gamma-ray spectra of a sample of blazars as a cutoff that changes amplitude and energy with redshift. Comparing the strength of the absorption feature to current EBL models yielded a EBL density of $3 (\pm 1) \text{ nW m}^{-2} \text{ sr}^{-1}$ at ultraviolet frequency for $z \approx 1$. This EBL level is close to lower limits expected from source counts and constrains the contribution of Active Galactic Nuclei and low-metallicity stars to the EBL in the early Universe.

Keywords: AGN, Extragalactic background Light, star formation, gamma-rays.

Information about the star formation history of the Universe is encoded in the intensity of the Extragalactic Background light (EBL). Direct detection of the EBL radiation fields is thus important, but at present extremely difficult [1]. An indirect means of probing the diffuse radiation fields is through γ - γ absorption of high-energy gamma rays [2, 3, 4]. In this process, a gamma-ray photon of energy E_γ and an EBL photon of energy E_{EBL} annihilate and create an electron-positron pair. This process occurs for head-on collisions when (e.g.) $E_\gamma \times E_{EBL} \geq 2(m_e c^2)^2$, where $m_e c^2$ is the rest mass energy of the electron. This introduces an attenuation in the spectra of gamma-ray sources above a critical gamma-ray energy of $E_{crit}(z) \approx 170(1+z)^{-2.38} \text{ GeV}$ [5].

Several attempts have been made to detect the long-sought EBL attenuation [7, 8, 9]. So far, limits on the EBL density have been inferred from the absence of absorption features in the spectra of individual blazars [10, 8], distant galaxies with bright gamma-ray emission powered by matter accreting onto central, massive black holes. While this feature is indeed difficult to constrain for a single source, we showed that it is detected collectively in the gamma-ray spectra of a sample of blazars as a cut-off that changes amplitude and energy with redshift. This result was published in [11] and will be described in the following. We note that the gamma-ray absorption feature has recently also been detected by other authors [12, 13], complementing the measurement presented here to lower EBL frequencies.

We searched for an attenuation of the spectra of blazars in the 1–500 GeV band using the first 46 months of observations of the Large Area Telescope (LAT) on board the *Fermi* satellite. At these energies gamma rays are absorbed by EBL photons in the optical to UV range. Due to the large energy and redshift coverage, *Fermi*-LAT

measures the intrinsic (i.e. unabsorbed) spectrum up to $\sim 100 \text{ GeV}$ for any blazar at $z < 0.2$, and up to $\sim 15 \text{ GeV}$ for any redshift. The data were filtered, removing time periods in which the instrument was not in sky-survey mode, and removing photons whose zenith angle is larger than 100° to limit contamination from the Earth limb emission. We consider only photons collected within 15° of the source position with $1 \leq E \leq 500 \text{ GeV}$. We employ the P7SOURCE_V6 instrumental response function (IRF) and perform a binned likelihood analysis. The Galactic and isotropic diffuse emissions are modeled using respectively the gal_2yearp7v6_v0.fits and iso_p7v6source.txt templates.

The LAT has detected > 1000 blazars to date [14]. We restricted our search to a subset of 150 blazars of the BL Lacertae (BL Lac) type that are significantly detected above 3 GeV, because of the expected lack of intrinsic absorption [15]. The sample covers a redshift range 0.03–1.6. The critical energy is therefore always $\geq 25 \text{ GeV}$, which means that the spectrum measured below this energy is unabsorbed and a true representation of the intrinsic spectrum of the source. We thus determined the intrinsic source spectrum relying on data between 1 GeV and the critical energy E_{crit} and extrapolated it to higher energies. By combining all the spectra we were able to determine, the average deviation, above the critical energy, of the measured spectra from the intrinsic ones, which provides a measurement of the optical depth $\tau_{\gamma\gamma}$.

We determined the spectral parameters of each blazar by maximizing the likelihood of a given source model. The model comprised the Galactic and isotropic diffuse components and all sources in the second *Fermi* LAT catalog [17] within a region of interest (ROI) of 15° radius. We modeled the spectra of the sources in our sample as parabolic in the logarithmic space of energy and flux (see Eq. 2 in [17])

for a definition). Their spectra were modified by a term $e^{-\tau_{\gamma\gamma}(E,z)}$ that describes the absorption of gamma-ray photons on the EBL. In the above we defined $\tau_{\gamma\gamma}(E,z) = b \cdot \tau_{\gamma\gamma}^{model}(E,z)$, where the $\tau_{\gamma\gamma}^{model}(E,z)$ is the optical depth predicted by EBL models [18, 19, 5, 20, 21] and b is a scaling variable, left free in the likelihood maximization. In particular, this allowed us to assess the likelihood of two scenarios: i) there is no EBL attenuation ($b=0$), ii) the model prediction is correct ($b=1$).

We combined the data from all the ROIs in a global fit that determined the common parameter b for a given EBL model [11]. All those models with a minimal EBL density based on (or compatible with) resolved galaxy counts [5, 6, 22, 20, 21, 23] were found to be acceptable descriptions of the *Fermi* data (i.e. are consistent with $b=1$ within $\approx 25\%$, see also Figure 2) yielding a significance of the absorption feature of up to $\sim 6\sigma$. Models that predict a larger intensity of the EBL particularly in the UV [18, 19] would produce a stronger-than-observed attenuation feature and are therefore incompatible with the *Fermi* observations. Our measurement points to a minimal level of the optical-UV EBL up to redshift $z \approx 1.6$ which combined with the upper limits [10, 24, 25] derived at lower redshift (using observations of blazars at TeV energies) on the near-infrared EBL highlights the conclusion that most of the EBL intensity can be explained by the measured galaxy emission.

Our measurement relies on the accuracy of the extrapolation of the intrinsic spectra of the sources above the critical energy [26]. This in turn depends on a precise description of the gamma-ray spectra by our source parametrization. To verify that this is the case and to exclude the possibility that the detected absorption feature is intrinsic to the gamma-ray sources [15], we performed the analysis in 3 independent redshift intervals ($z < 0.2$, $0.2 \leq z < 0.5$, and $0.5 \leq z < 1.6$). The deviations from the intrinsic spectra in the three redshift intervals are displayed in Figure 1. In the local Universe ($z < 0.2$), EBL absorption is negligible in most of the *Fermi*-LAT energy band ($E_{crit} \geq 120$ GeV). The lowest redshift interval therefore shows directly the intrinsic spectra of the sources and shows that our spectral parametrization is accurate. The absorption feature is clearly visible above the critical energy in the higher redshift bins. Its amplitude and modulation in energy evolve with redshift as expected for EBL absorption. In principle, the observed attenuation could be due to a spectral cut-off that is intrinsic to the gamma-ray sources. The absence of a cut-off in the spectra of sources with $z < 0.2$ would require that the properties of BL Lacs change with redshift or luminosity. It remains an issue of debate whether such evolution exists [27, 28, 29, 30]. However, in case it were present, the intrinsic cut-off would be expected to evolve differently with redshift than we observe. To illustrate this effect, we fitted the blazar sample assuming that all the sources have an exponential cut-off at an energy E_0 . From source to source the observed cut-off energy changes because of the source redshift and because we assumed that blazars as a population are distributed in a sequence such as that proposed in [27, 28, 29, 30]. E_0 was fitted to the data globally like b above. As apparent from Figure 1, it appears difficult to reconcile the observed feature with an intrinsic characteristic of the blazars' spectra. We therefore associate the spectral feature to the EBL absorption.

At energies ≤ 100 GeV, gamma rays observed at Earth and coming from redshift ≥ 1 interact mostly with UV photons of ≥ 5 electron volts. An UV background in excess

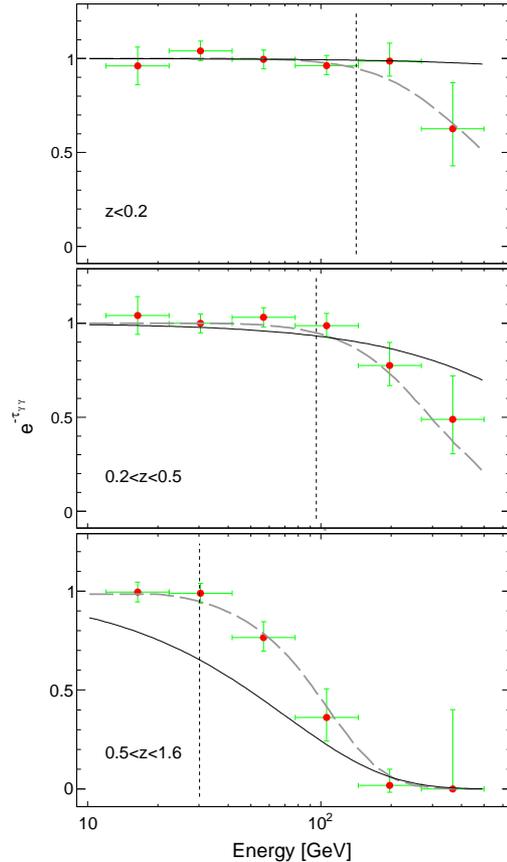


Figure 1: Absorption feature present in the spectra of BL Lacertae objects as a function of increasing redshift (data points, from top to bottom). The dashed curves show the attenuation expected for the sample of sources by averaging, in each redshift and energy bin, the opacities of the sample (the model of [5] was used) and multiplying this average by the best-fit scaling parameter b obtained independently in each redshift interval. The vertical line shows the critical energy E_{crit} below which $\leq 5\%$ of the source photons are absorbed by the EBL. The thin solid curve represents the best-fit model assuming that all the sources have an intrinsic exponential cut-off and that blazars follow the blazar sequence model of [28, 29]. More details on the model are found in the online supplements of [11].

of the light emitted by resolved galaxies can be produced locally by AGN or at higher redshift ($z \approx 7-15$) by low-metallicity massive stars [31]. By comparing the results from the best-fit EBL models, we measured the UV component of the EBL to have an intensity of $3(\pm 1) \text{ nW m}^{-2} \text{ sr}^{-1}$ at $z \approx 1$. A contribution to the UV background from AGN as large as the one predicted by [32] (i.e., $\approx 10 \text{ nW m}^{-2} \text{ sr}^{-1}$) and used in the EBL model of [18] is thus excluded by our analysis at high confidence. However, the recent prediction [33] of the UV background from AGN ($\approx 2 \text{ nW m}^{-2} \text{ sr}^{-1}$) is in agreement with the *Fermi* measurement. Direct measurements of the extragalactic UV background are hampered by the strong dust-scattered Galactic radiation [34]. The agreement between the intensity of the UV background as measured with *Fermi* and that due to galaxies individually resolved by the Hubble Space Telescope

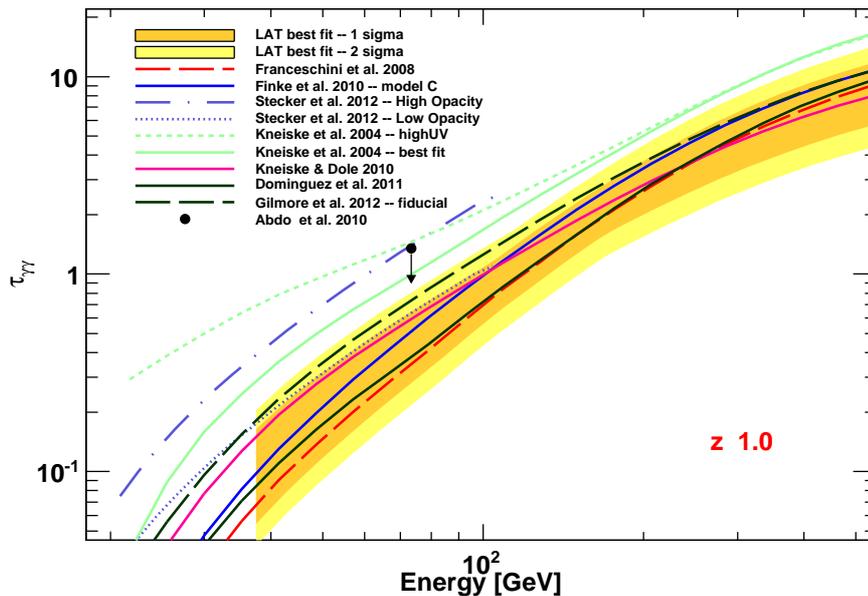


Figure 2: Measurement, at the 68 % and 95 % confidence levels (including systematic uncertainties added in quadrature), of the opacity $\tau_{\gamma\gamma}$ from the best fits to the *Fermi* data compared to predictions of EBL models. The plot shows the measurement at $z \approx 1$ which is the average redshift of the most constraining redshift interval (i.e. $0.5 \leq z < 1.6$). The *Fermi*-LAT measurement was derived combining the limits on the best-fit EBL models. The downward arrow represents the 95 % upper limit on the opacity at $z=1.05$ derived in [8]. For clarity this figure shows only a selection of the models we tested while the full list is reported in the online supplements of [11]. The EBL models of [45], which are not defined for $E \geq 250/(1+z)$ GeV and thus could not be used, are reported here for completeness.

[35] ($3 \pm 1 \text{ nW m}^{-2} \text{ sr}^{-1}$ versus $2.9\text{--}3.9 \text{ nW m}^{-2} \text{ sr}^{-1}$, respectively) shows that the room for any residual diffuse UV emission is small. This conclusion is reinforced by the good agreement of the *Fermi* measurement and the estimate of the average UV background, at $z \geq 1.7$, of $2.2\text{--}4.0 \text{ nW m}^{-2} \text{ sr}^{-1}$ using the proximity effect in quasar spectra [36].

Zero-metallicity population-III stars or low-metallicity population-II stars are thought to be the first stars to form in the Universe and formally marked the end of the dark ages when, with their UV light, these objects started ionizing the intergalactic medium [37]. These stars are also believed to be responsible for creating the first metals and dispersing them in the intergalactic medium [38, 39, 40]. A very large contribution of population-III stars to the near-infrared EBL had already been excluded by [10]. Our measurement constrains, according to [41, 42], the redshift of maximum formation of low-metallicity stars to be at $z \geq 10$ and its peak co-moving star-formation rate to be lower than $0.5 M_{\odot} \text{ Mpc}^{-3} \text{ yr}^{-1}$. This upper limit is of the same order of the peak star-formation rate of $0.2\text{--}0.6 M_{\odot} \text{ Mpc}^{-3} \text{ yr}^{-1}$ proposed by [43] and suggests that the peak star-formation rate might be much lower as proposed by [44].

In the future we expect to increase the accuracy and redshift coverage of our measurement due to two principal reasons: First, the number of LAT-detected BL Lacs with measured redshift is continuously increasing [65]. Second we expect our measurement to be applicable to the Blazar class of Flat Spectrum Radio Quasars. The latter are as numerous as BL Lacs in the gamma-ray sky and are typically located at higher redshifts.

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References

- [1] S. Inoue, *et al.*, *Monthly Notices of the Royal Astronomical Society* **404**, 1938 (2010).
- [2] R. J. Gould, G. Schröder, *Physical Review Letters* **16**, 252 (1966).
- [3] G. G. Fazio, F. W. Stecker, *Nature* **226**, 135 (1970).
- [4] F. W. Stecker, O. C. de Jager, M. H. Salamon, *The Astrophysical Journal* **390**, L49 (1992).
- [5] A. Franceschini, G. Rodighiero, M. Vaccari, *Astronomy and Astrophysics* **487**, 837 (2008).
- [6] R. C. Gilmore, P. Madau, J. R. Primack, R. S. Somerville, F. Haardt, *Monthly Notices of the Royal Astronomical Society* **399**, 1694 (2009).
- [7] K. Mannheim, S. Westerhoff, H. Meyer, H. H. Fink, *The Astrophysical Journal* **315**, 77 (1996).

- [8] A. A. Abdo, *et al.*, *The Astrophysical Journal* **723**, 1082 (2010).
- [9] M. Raue, *Astronomy and Astrophysics* **520**, A34 (2010).
- [10] F. Aharonian, *et al.*, *Nature* **440**, 1018 (2006).
- [11] M. Ackermann, *et al.*, *Science* **338**, 1190 (2012).
- [12] A. Domínguez *et al.*, *The Astrophysical Journal* **770**, 77 (2013).
- [13] A. Abramowski *et al.*, *Astronomy and Astrophysics* **550**, A4 (2013).
- [14] M. Ackermann, *et al.*, *The Astrophysical Journal* **743**, 171 (2011).
- [15] A. Reimer, *The Astrophysical Journal* **665**, 1023 (2007).
- [16] The software to perform analysis of *Fermi* data is available at <http://fermi.gsfc.nasa.gov/ssc/data/analysis/software/>.
- [17] P. Nolan, *et al.*, *arXiv:1108.1435* (2011).
- [18] T. M. Kneiske, T. Bretz, K. Mannheim, D. H. Hartmann, *Astronomy and Astrophysics* **413**, 807 (2004).
- [19] F. W. Stecker, M. A. Malkan, S. T. Scully, *The Astrophysical Journal* **648**, 774 (2006).
- [20] J. D. Finke, S. Razzaque, C. D. Dermer, *The Astrophysical Journal* **712**, 238 (2010).
- [21] A. Domínguez, *et al.*, *Monthly Notices of the Royal Astronomical Society* **410**, 2556 (2011).
- [22] T. M. Kneiske, H. Dole, *Astronomy and Astrophysics* **515**, A19 (2010).
- [23] R. C. Gilmore, R. S. Somerville, J. R. Primack, A. Domínguez, *Monthly Notices of the Royal Astronomical Society* **422**, 3189 (2012).
- [24] D. Mazin, M. Raue, *Astronomy and Astrophysics* **471**, 439 (2007).
- [25] The MAGIC Collaboration, *et al.*, *Science* **320**, 1752 (2008).
- [26] We show in the supplemental material that our result is robust against conservative choices the critical energy.
- [27] G. Ghisellini, A. Celotti, G. Fossati, L. Maraschi, A. Comastri, *Monthly Notices of the Royal Astronomical Society* **301**, 451 (1998).
- [28] G. Fossati, L. Maraschi, A. Celotti, A. Comastri, G. Ghisellini, *Monthly Notices of the Royal Astronomical Society* **299**, 433 (1998).
- [29] G. Ghisellini, L. Maraschi, F. Tavecchio, *Monthly Notices of the Royal Astronomical Society* **396**, L105 (2009).
- [30] E. T. Meyer, G. Fossati, M. Georganopoulos, M. L. Lister, *The Astrophysical Journal* **740**, 98 (2011).
- [31] M. R. Santos, V. Bromm, M. Kamionkowski, *Monthly Notices of the Royal Astronomical Society* **336**, 1082 (2002).
- [32] F. Haardt, P. Madau, *The Astrophysical Journal* **461**, 20 (1996).
- [33] F. Haardt, P. Madau, *The Astrophysical Journal* **746**, 125 (2012).
- [34] S. Bowyer, *Annual Review of Astronomy and Astrophysics* **29**, 59 (1991).
- [35] J. P. Gardner, T. M. Brown, H. C. Ferguson, *The Astrophysical Journal* **542**, L79 (2000).
- [36] J. Scott, J. Bechtold, A. Dobrzycki, V. P. Kulkarni, *The Astrophysical Journal* **130**, 67 (2000).
- [37] V. Bromm, R. B. Larson, *Annual Review of Astronomy and Astrophysics* **42**, 79 (2004).
- [38] J. P. Ostriker, N. Y. Gnedin, *The Astrophysical Journal* **472**, L63 (1996).
- [39] T. H. Greif, J. L. Johnson, V. Bromm, R. S. Klessen, *The Astrophysical Journal* **670**, 1 (2007).
- [40] J. H. Wise, T. Abel, *The Astrophysical Journal* **685**, 40 (2008).
- [41] M. Raue, T. Kneiske, D. Mazin, *Astronomy and Astrophysics* **498**, 25 (2009).
- [42] R. C. Gilmore, *eprint arXiv:1109.0592* (2011).
- [43] V. Bromm, A. Loeb, *The Astrophysical Journal* **575**, 111 (2002).
- [44] L. Tornatore, A. Ferrara, R. Schneider, *Monthly Notices of the Royal Astronomical Society* **382**, 945 (2007).
- [45] F. W. Stecker, M. A. Malkan, S. T. Scully, *arXiv:1205.5168* (2012).
- [46] J. T. Stocke, *et al.*, *The Astrophysical Journals* **76**, 813 (1991).
- [47] C. M. Urry, P. Padovani, *Publications of the Astronomical Society of the Pacific* **107**, 803 (1995).
- [48] M. J. M. Marcha, I. W. A. Browne, C. D. Impey, P. S. Smith, *Monthly Notices of the Royal Astronomical Society* **281**, 425 (1996).
- [49] A. A. Abdo, *et al.*, *The Astrophysical Journal* **736**, 131 (2011).
- [50] A. A. Abdo, *et al.*, *The Astrophysical Journal* **727**, 129 (2011).
- [51] E. Aliu, *et al.*, *The Astrophysical Journal* **750**, 94 (2012).
- [52] J. R. Primack, J. S. Bullock, R. S. Somerville, *High Energy Gamma-Ray Astronomy*, F. A. Aharonian, H. J. Völk, & D. Horns, ed. (2005), vol. 745 of *American Institute of Physics Conference Series*, pp. 23–33.
- [53] A. Tramacere, E. Massaro, A. M. Taylor, *The Astrophysical Journal* **739**, 66 (2011).
- [54] E. Massaro, A. Tramacere, M. Perri, P. Giommi, G. Tosti, *Astronomy and Astrophysics* **448**, 861 (2006).
- [55] Fermi-LAT Collaboration, *arXiv:1206.1896* (2012).
- [56] R. Plaga, *Nature* **374**, 430 (1995).
- [57] F. Tavecchio, G. Ghisellini, G. Bonnoli, L. Foschini, *Monthly Notices of the Royal Astronomical Society* **414**, 3566 (2011).
- [58] A. Neronov, D. Semikoz, M. Kachelriess, S. Ostapchenko, A. Elyiv, *The Astrophysical Journal* **719**, L130 (2010).
- [59] I. Vovk, A. M. Taylor, D. Semikoz, A. Neronov, *The Astrophysical Journal* **747**, L14 (2012).
- [60] C. D. Dermer, *et al.*, *The Astrophysical Journal* **733**, L21 (2011).
- [61] F. Tavecchio, G. Ghisellini, G. Ghirlanda, L. Foschini, L. Maraschi, *Monthly Notices of the Royal Astronomical Society* **401**, 1570 (2010).
- [62] K. Mannheim, *The Astrophysical Journal* **269**, 67 (1993).
- [63] E. Waxman, P. Coppi, *The Astrophysical Journal* **464**, L75 (1996).
- [64] W. Essey, A. Kusenko, *Astroparticle Physics* **33**, 81 (2010).
- [65] M. S. Shaw, *et al.*, *ArXiv:1301.0323* (2013)